

## Error threshold in the evolution of diploid organisms

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

1997 J. Phys. A: Math. Gen. 30 2601

(<http://iopscience.iop.org/0305-4470/30/8/009>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 171.66.16.112

The article was downloaded on 02/06/2010 at 06:16

Please note that [terms and conditions apply](#).

## Error threshold in the evolution of diploid organisms

D Alves and J F Fontanari

Instituto de Física de São Carlos, Universidade de São Paulo, Caixa Postal 369, 13560-970 São Carlos SP, Brazil

Received 26 November 1996

**Abstract.** The effects of error propagation in the reproduction of diploid organisms are studied within the population genetics framework of the quasispecies model. The dependence of the error threshold on the dominance parameter is fully investigated. In particular, it is shown that dominance can protect the wild-type alleles from the error catastrophe. The analysis is restricted to a diploid analogue of the single-peaked fitness landscape.

The finding that the length of self-reproducing molecules that compete for a finite supply of resources is limited by their replication accuracy is probably the main outcome of Eigen's quasispecies model (Eigen 1971). This phenomenon, termed error threshold, poses an interesting challenge to the theories of the origin of life, since it prevents the emergence of huge molecules which could carry the necessary information for building a complex metabolism (Eigen and Schuster 1979, Kauffman 1993).

In the quasispecies model, a molecule is represented by a string of  $\nu$  digits  $s = (s_1, s_2, \dots, s_\nu)$ , with the variables  $s_i$  allowed to take on  $\kappa$  different values, each representing a different type of monomer used to build the molecule. The focus is on the time evolution of the concentrations  $x_i$  of molecules of type  $i = 1, 2, \dots, \kappa^\nu$  which obey the following differential equations (Eigen 1971):

$$\frac{dx_i}{dt} = \sum_j W_{ij}x_j - [D_i + \Phi(t)]x_i \quad (1)$$

where the constants  $D_i$  stand for the death probability of molecules of type  $i$  and  $\Phi(t)$  is a dilution flux that keeps the total concentration constant. The feature that distinguishes this model from the well established models of population genetics (Hartl and Clark 1989) is the replication matrix  $W$  which takes into account the primary structure of the molecules. More specifically, its elements are given by

$$W_{ii} = A_i q^\nu \quad (2)$$

and

$$W_{ij} = \frac{A_i}{(\kappa - 1)^{d(i,j)}} q^{\nu - d(i,j)} (1 - q)^{d(i,j)} \quad i \neq j \quad (3)$$

where  $A_i$  is the replication rate of molecules of type  $i$ ,  $d(i, j)$  is the Hamming distance between molecules  $i$  and  $j$ , and  $q \in [0, 1]$  is the single-monomer replication accuracy, which is assumed to be the same for all monomers.

For simple replication landscapes, the solutions of the  $\kappa^v$  kinetic equations (1) have been thoroughly studied using perturbation theory (Eigen *et al* 1989). More complex, spin-glass-like replication landscapes can be analysed using the correspondence between those equations and the equilibrium properties of a surface lattice system (Leuthäusser 1986, 1987, Tarazona 1992, Franz *et al* 1993). It is worth mentioning that for the single-peaked replication landscape the exact stationary solution of equations (1) can be obtained by mapping them into a polymer localization problem (Galluccio *et al* 1996). Recently, a population genetics approach to the quasispecies model has been proposed that, in spite of its simplicity, yields results that are qualitatively similar to those obtained by solving the kinetic equations (Alves and Fontanari 1996).

An alternative interpretation of the quasispecies model is given by considering the  $\kappa^v$  different strings  $s$  as different forms (alleles) of a certain gene that determines the fitness  $A_i$  of the haploid organisms. Thus, this model is equivalent to the classical one-locus, multiple-allele model of population genetics (Hartl and Clark 1989), except for the mutation mechanism which must be adapted to satisfy the constraints imposed by the internal structure of the alleles. Accordingly, Wiehe *et al* (1995) have generalized the original haploid formulation of the quasispecies model so as to consider the evolution of diploid organisms as well. An important by-product of that analysis is the study of the effects of dominance on error thresholds, which has led to an interesting conjecture about the evolution of dominance.

In this paper we employ the population genetics formulation of the quasispecies model to investigate the error propagation in the reproduction of diploid organisms. This approach allows us to study in great detail the dynamical behaviour of the model in the full space of the control parameters  $v$  and  $q$  as well as in the space of the parameters that specify the fitness landscapes. We should mention that the analysis of Wiehe *et al* (1995) was based on the numerical solution of the diploid counterpart of the kinetic equations (1) and on very crude approximations that neglect the effects of back mutations.

In the population genetics formulation, the  $\kappa^v$  different alleles are grouped into  $(v+\kappa-1)!/v!(\kappa-1)!$  classes, according to the number of monomers of each type they have, regardless of their specific position inside the allele. Hence, a given class is characterized by the vector  $\mathbf{P} = (P_1, P_2, \dots, P_\kappa)$ , where  $P_\alpha$  is the number of monomers of type  $\alpha$  in any allele inside that class. Clearly,  $\sum_\alpha P_\alpha = v$ . The alleles belonging to the same class are assumed to be equivalent, in the sense that their presence confers the same fitness value on the genotypes. The crucial simplifying assumption of the population genetics approach is that, given the monomer frequencies in generation  $t$ ,  $p_\alpha(t)$  with  $\sum_\alpha p_\alpha(t) = 1$ , the frequencies of alleles in class  $\mathbf{P}$  are given by the multinomial distribution

$$\Pi_t(\mathbf{P}) = C_{\mathbf{P}}^v [p_1(t)]^{P_1} [p_2(t)]^{P_2} \dots [p_\kappa(t)]^{P_\kappa} \quad (4)$$

where  $C_{\mathbf{P}}^v = v!/P_1!P_2!\dots P_\kappa!$ . Thus, at generation  $t$ , the monomers are sampled with replacement from an urn containing  $\kappa$  different types of monomers in the proportions  $p_\alpha(t)$ ;  $\alpha = 1, \dots, \kappa$ .

Let  $A(\mathbf{P}^i, \mathbf{P}^j) = A(\mathbf{P}^j, \mathbf{P}^i)$  denote the fitness of the genotypes  $\mathbf{P}^i \mathbf{P}^j$ , i.e. genotypes composed of any pair of alleles belonging to classes  $\mathbf{P}^i$  and  $\mathbf{P}^j$ . Then the fraction of monomers  $\alpha$  that the genotype  $\mathbf{P}^i \mathbf{P}^j$  contributes to generation  $t+1$  is proportional to the product of three factors: (a) its frequency in the population  $\Pi_t(\mathbf{P}^i, \mathbf{P}^j)$ , (b) its fitness  $A(\mathbf{P}^i, \mathbf{P}^j)$ , and (c) the average number of monomers  $\alpha$  that replicate correctly,  $q(P_\alpha^i + P_\alpha^j)$ , plus the average number of monomers  $\beta \neq \alpha$  that mutate to  $\alpha$ ,  $[(1-q)/(\kappa-1)] \sum_{\beta \neq \alpha} (P_\beta^i + P_\beta^j)$ . A simple calculation yields the following equations for the time evolution of the

monomer frequencies:

$$p_\alpha(t + 1) = \frac{1}{\kappa - 1} \left[ 1 - q + \frac{\kappa q - 1}{2w_t} \sum_{P^i} \sum_{P^j} \Pi_t(P^i, P^j) A(P^i, P^j) (P_\alpha^i + P_\alpha^j) \right] \tag{5}$$

where

$$\Pi_t(P^i, P^j) = \Pi_t(P^i) \Pi_t(P^j) \tag{6}$$

and the normalization factor,

$$w_t = \nu \sum_{P^i} \sum_{P^j} \Pi_t(P^i, P^j) A(P^i, P^j) \tag{7}$$

is the average fitness of the entire population. Here the notation  $\sum_P$  stands for  $\sum_{P_1=0}^\nu \dots \sum_{P_\kappa=0}^\nu \delta(\nu, \sum_\alpha P_\alpha)$ , where  $\delta(k, l)$  is the Kronecker delta. It is interesting to note that equation (5) is identical to the equation governing the evolution of sexually reproducing haploid organisms (Alves and Fontanari 1996).

In the remainder of this paper we will consider binary strings only. In this case there are two types of monomers ( $\kappa = 2$ ), so that the alleles are characterized by a single parameter, namely, the number of monomers of type 1 they have,  $P_1 \equiv P$ . The extension of our analysis to larger values of  $\kappa$  is straightforward. To proceed further we must specify the fitness of the genotypes  $P^i P^j$ . According to Wiehe *et al* (1995) we consider the following diploid analogue of the single-peaked fitness landscape:

$$A(P^i, P^j) = \begin{cases} (1 + a)^2 & \text{if } P^i = P^j = \nu \\ (1 + a)^{2h} & \text{if } P^i = \nu \text{ and } P^j \neq \nu \\ 1 & \text{if } P^i \neq \nu \text{ and } P^j \neq \nu \end{cases} \tag{8}$$

where  $a > 0$  is the parameter measuring the selective advantage of the so-called *master* allele  $P = \nu$ , and  $-\infty < h < \infty$  is the dominance parameter. The master allele is completely dominant for  $h = 1$  and completely recessive for  $h = 0$ . For  $h = \frac{1}{2}$  we find  $A(P^i, P^j) = A(P^i)A(P^j)$  and so there is no dominance. In this case equation (5) reduces to the equation that governs the evolution of asexually reproducing haploid organisms (Alves and Fontanari 1996). Thus the intervals  $h \in [0, \frac{1}{2})$  and  $h \in (\frac{1}{2}, 1]$  delimit the regions of recessivity and dominance, respectively, of the master allele. There are other cases of interest as well:  $h > 1$  models the phenomenon of heterosis or hybrid vigour (heterozygote advantage), while  $h < 0$  models the phenomenon that occurs at the early stages of speciation when hybrids are less viable (heterozygote disadvantage).

Inserting equation (8) into the recurrence equation (5) yields the following equation for the frequency of monomers of type 1 in generation  $t$ ,  $p_1(t) \equiv p_t$ :

$$p_{t+1} = 1 - q + (2q - 1) \frac{\Lambda_1 p_t^{2\nu} + \Lambda_2 (p_t + 1) p_t^\nu + p_t}{\Lambda_1 p_t^{2\nu} + 2\Lambda_2 p_t^\nu + 1} \tag{9}$$

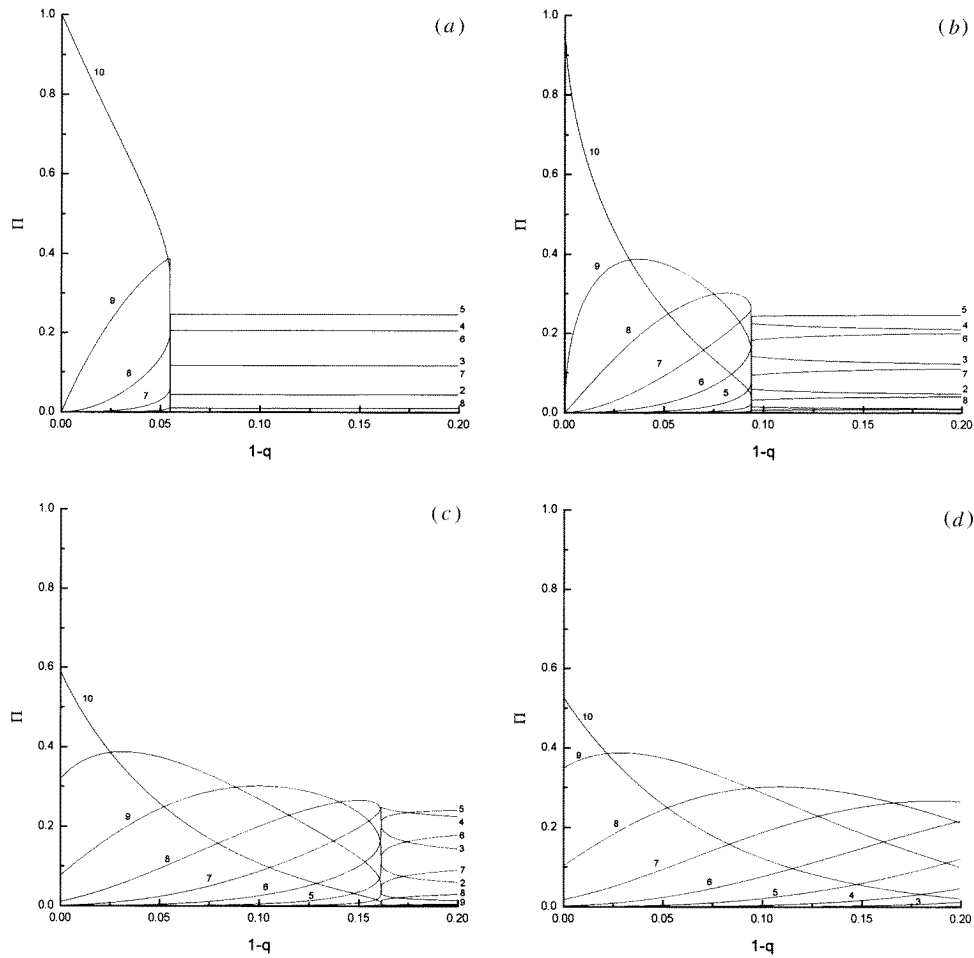
where

$$\Lambda_1 = (1 + a)^2 - 2(1 + a)^{2h} + 1 \tag{10}$$

and

$$\Lambda_2 = (1 + a)^{2h} - 1. \tag{11}$$

In figure 1 we present the steady-state frequencies of alleles, obtained by solving the recursion equation (9) with  $p_0 \approx 1$ , as a function of the error rate per monomer  $1 - q$  for  $\nu = 10$ ,  $a = 2$ , and different values of the dominance parameter. In the case of perfect replication accuracy ( $1 - q = 0$ ), the fixed point  $p^* = 0$  is always unstable, while  $p^* = 1$

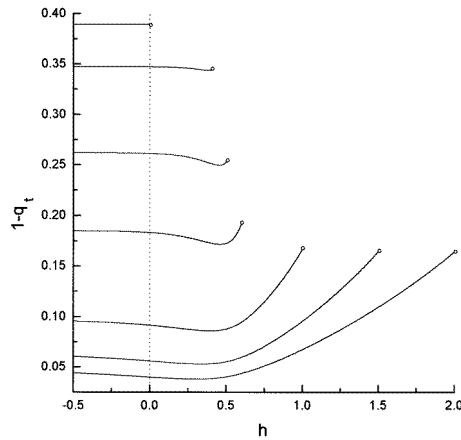


**Figure 1.** Steady-state frequencies of alleles belonging to classes  $P = 10$  (master allele) to  $P = 0$  as a function of the error rate per digit  $1 - q$  for  $\nu = 10$ ,  $a = 2$ , and (a)  $h = 0$ , (b)  $h = 1$ , (c)  $h = 1.5$ , and (d)  $h = 2$ .

is stable for  $h \leq 1$  only. For  $h > 1$ , a third (stable) fixed point  $\frac{1}{2} < p^* \approx 1$  appears, signalling the emergence of heterosis. For  $h \leq h_c \approx 1.75$  there are two distinct regimes: the *quasispecies* regime characterized by a population dominated by the master allele and its close neighbours, and the *uniform* regime where the  $2^\nu$  alleles appear in the same proportion (clearly the class  $P = \nu/2$  is the most favoured in this case). The error rate at which the discontinuous transition between these two regimes takes place is termed *error threshold*  $1 - q_t$ . As  $h$  increases, the size of the jump at the transition decreases till it disappears at a critical value  $h = h_c$ . Beyond that value it is no longer possible to distinguish the two regimes.

To better characterize the error threshold transition we concentrate our analysis on the nature of the fixed points  $p_{t+1} = p_t = p^*$  which are given by the real roots of  $f(p) = 0$ , where

$$f(p) = \Lambda_1(p - q)p^{2\nu} + \Lambda_2(3p - 2pq - 1)p^\nu - (1 - q)(1 - 2p). \quad (12)$$

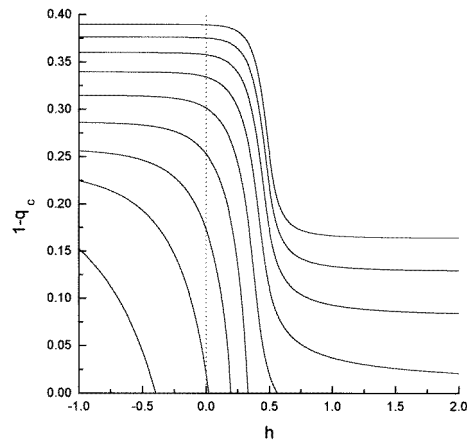


**Figure 2.** Error threshold  $1 - q_t$  as a function of the dominance parameter  $h$  for  $\nu = 10$  and (from top to bottom)  $a = 314.8, 186.1, 57.0, 18.8, 4.4, 2.0,$  and  $1.3$ . The parameter  $a$  was chosen so that the transition lines end at critical points located at  $h = 0, 0.4, 0.5, 0.6, 1.0, 1.5$  and  $2.0$ , respectively.

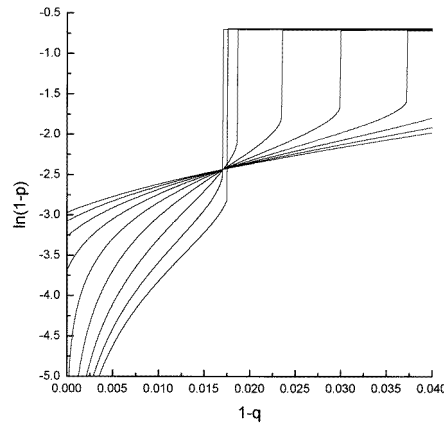
For small error rates this equation has only one real root which corresponds to the stable fixed point  $p^* \approx 1$  associated to the quasispecies regime. As the error rate increases, a double root appears, originating two new fixed points: a stable one,  $p^* \approx \frac{1}{2}$ , associated to the uniform regime, and an unstable one that delimitates the basins of attraction of the stable fixed points. These fixed points co-exist till the error rate reaches the threshold value  $1 - q_t$ , where the stable quasispecies fixed point and the unstable one coalesce. For larger error rates, equation (12) has only one real root which corresponds to the uniform fixed point. Thus, the error threshold transition can be easily determined by solving  $f(p) = df(p)/dp = 0$  simultaneously for  $p$  and  $q = q_t$ . As mentioned above, since these equations have two solutions we must choose the one with the larger value of  $p$ . The critical point  $h_c$  is determined by tuning the value of  $h$  so that the three real roots of (12) coincide, i.e. we have to solve the three equations  $f(p) = df(p)/dp = d^2f(p)/dp^2 = 0$  simultaneously for  $p, q = q_c$  and  $h = h_c$ .

Using the prescriptions given above, we present in figure 2 the error threshold transition lines as a function of  $h$  for  $\nu = 10$  and several values of  $a$ . The error threshold  $1 - q_t$  is practically insensitive to variations of  $h$  for negative and small positive values of this parameter. It reaches its minimal value around  $h = 0.5$  (non-dominance regime) and then increases quickly as the system enters the dominance region,  $h > 0.5$ . We note the re-entrant behaviour of these transition lines: for certain values of  $q$ , the system undergoes two discontinuous transitions as  $h$  is increased. The transition lines end at critical points, which are shown in figure 3 for different values of  $\nu$ . It is interesting to note that only for  $\nu < 7$  (or more exactly  $\nu < 6.93$ ) the critical lines touch the axis  $1 - q_c = 0$ . So, in these cases, there are values of the dominance parameter  $h > 0$  for which the error threshold transition never occurs.

Until now we have concentrated on the location of the error threshold as a function of the dominance parameter. We turn now to the analysis of the composition of the population at the steady state. It can be characterized by the average normalized Hamming distance from the master allele which, within the population genetics framework, is given simply by  $1 - p^*$ . This quantity is shown in figure 4 as a function of the error rate for  $\nu = 10$ ,



**Figure 3.** Error threshold at the critical point  $1 - q_c$  as a function of the dominance parameter  $h$  for (from top to bottom)  $\nu = 10$  to  $\nu = 2$ .



**Figure 4.** Average normalized Hamming distance from the master as a function of the replication error rate for  $\nu = 10$ ,  $a = 0.5$ , and (from top to bottom before the intersection)  $h = 2, 1.75, 1.5, 1.25, 1, 0.75, 0.5, 0.25$ , and  $0$ . The curves for  $h > h_s = 0.244$  intersect at the inflection point  $1 - q_r = 0.017$ .

$a = 0.5$  and several values of  $h$ . What is remarkable about this figure is that there exists a value of the error rate  $1 - q = 1 - q_r \approx 0.017$  such that the fixed point  $p^* \approx 0.912$  is independent of  $h$ . This fixed point, however, becomes unstable for  $h < h_s \approx 0.244$ . Thus, although recessivity leads to a higher concentration of the master allele for  $1 - q < 1 - q_r$ , this allele is quickly lost from the population for larger error rates. The main effect of dominance is to postpone the error catastrophe at the price of reducing the concentration of the master allele in the population. We note that at the inflection point  $q = q_r$  the effects of dominance and recessivity are reversed. This point can be easily determined by setting to zero the coefficient of the term  $(1 + a)^{2h}$  in equation (12), namely,

$$g(p) = -2(p - q)p^\nu + (3 - 2q)p - 1 \quad (13)$$

and solving  $g(p) = 0$  together with  $f(p) = 0$  for  $p$  and  $q = q_r$ .

Both analyses, the location of the error threshold and the composition of the population, indicate that dominance allows the master allele to resist to higher replication error rates than in the case of non-dominance. Actually, for sufficiently large  $h$ , it can even avoid the error catastrophe. This finding has been proposed as a possible explanation for the fact that the wild type, i.e. the allele that predominates in a population and that is particularly well suited to its environment, is often dominant: the dominant alleles might be the prevailing wild-type ones simply because they can tolerate higher error rates (Wiehe *et al* 1995).

In summary, we have employed the population genetics approach to the quasispecies model to investigate the error threshold catastrophe in the evolution of diploid organisms. In order to enhance the non-trivial effects of the imperfect replication accuracy of the organisms on the population composition, we have focused on a simple diploid analogue of the single-peaked fitness landscape. Two distinct steady-state regimes are observed: the quasispecies regime where the information about the environment, modelled by the fitness landscape, is preserved in the population composition, and the uniform regime, where this information is irreversibly lost. In the space of the parameters  $1 - q$  and  $h$ , these regimes are separated by discontinuous transition lines that terminate at critical points, beyond which

they become indistinguishable. We have found that dominance ( $h > 0.5$ ) can postpone or even avoid the error catastrophe. It is interesting, however, that a recessive allele ( $h < 0.5$ ) can do better than a non-dominant one ( $h \approx 0.5$ ).

To conclude, we mention that our results are in qualitative agreement with those of Wiehe *et al* (1995). Since the population genetics approach of the quasispecies model incorporates only a few essential features of the original chemical kinetics formulation, this agreement gives strong evidence for the robustness of the main conclusion drawn from the model, namely, the existence of an error catastrophe that limits the size of self-replicating organisms. Rather than just a caricature of the original model, the population genetics model presented in this paper may be viewed as a simpler, alternative model for investigating the evolution of self-replicating organisms, which may greatly facilitate the analysis of difficult problems such as the error propagation in finite populations and the effects of cooperation or catalysis among the evolving organisms.

### Acknowledgment

This work was supported in part by Conselho Nacional de Desenvolvimento Científico e Tecnológico (CNPq).

### References

- Alves D and Fontanari J F 1996 *Phys. Rev. E* **54** 4048  
Eigen M 1971 *Naturwissenschaften* **58** 465  
Eigen M, McCaskill J and Schuster P 1989 *Adv. Chem. Phys.* **75** 149  
Eigen M and Schuster P 1979 *The Hypercycle—A Principle of Natural Self-organization* (Berlin: Springer)  
Franz S, Peliti L and Sellitto M 1993 *J. Phys. A: Math. Gen.* **26** L1195  
Galluccio S, Graber R and Zhang Y-C 1996 *J. Phys. A: Math. Gen.* **29** L249  
Hartl D L and Clark A G 1989 *Principles of Population Genetics* (Sunderland: Sinauer Associates)  
Kauffman S A 1993 *The Origins of Order* (Oxford: Oxford University Press)  
Leuthäusser I 1986 *J. Chem. Phys.* **84** 1884  
———1987 *J. Stat. Phys.* **48** 343  
Tarazona P 1992 *Phys. Rev. A* **45** 6038  
Wiehe T, Baake E and Schuster P 1995 *J. Theor. Biol.* **177** 1